Homework 13. Due October 26

Problem 1. 3x5 points. Beam envelope in straight section.

For a one-dimensional motion consider beam propagating in a straight section starting as s_o and having length L. Let's eigen vector (beam envelope) at s_o is given by:

$$Y(\mathrm{s}_o) = \begin{bmatrix} \mathrm{w}_o \\ \mathrm{w}'_o + \dfrac{i}{\mathrm{w}_o} \end{bmatrix}; \beta_o \equiv \mathrm{w}_o^2;\ \ \alpha_o = -\frac{\beta'}{2} \equiv -\mathrm{w}_o \mathrm{w}'_o; \tag{1}$$

(a) Propagate the eigen vector along the straight section. Show that β-function can be expressed as

$$\beta(\mathrm{s}) = \beta^* + \frac{(s-s^*)^2}{\beta^*};$$

where β^*, s^* can be found from initial conditions (1). Hint, use derivative of β-function to find s^*. β^* is frequently used in colliders to describe the beam envelope in detectors.

(b) Calculate the (betatron) phase advance acquired in the straight section. Express it using β^*, s^*. Write expression for x(s) and x'(s). Show that x'=const.

(c) What is the maximum possible phase advance in a straight section (e.g. when $\mathrm{s_o}$,L are unlimited)?

Solution: (a) Propagating the eigen vector through a drift is just multiplying it by the drifts transport matrix:

$$\tilde{Y}(\mathrm{s}) = \begin{bmatrix} 1 & \Delta\mathrm{s} \\ 0 & 1 \end{bmatrix} Y(\mathrm{s}_o) = \begin{bmatrix} \mathrm{w}_o + \Delta\mathrm{s}\left(\mathrm{w}'_o + \dfrac{i}{\mathrm{w}_o}\right) \\ \mathrm{w}'_o + \dfrac{i}{\mathrm{w}_o} \end{bmatrix} = \begin{bmatrix} \mathrm{w}(\mathrm{s}) \\ \mathrm{w}'(\mathrm{s}) + \dfrac{i}{\mathrm{w}(\mathrm{s})} \end{bmatrix} e^{i\Delta\psi}; \Delta\mathrm{s} = \mathrm{s} - \mathrm{s_o}; \tag{2}$$

$$\beta(\mathrm{s}) = \mathrm{w}^2(\mathrm{s}) = \left|\mathrm{w}_o + \Delta\mathrm{s}\left(\mathrm{w}'_o + \frac{i}{\mathrm{w}_o}\right)\right|^2 = \left(\mathrm{w}_o + \Delta\mathrm{s}\mathrm{w}'_o\right)^2 + \frac{\Delta\mathrm{s}^2}{\mathrm{w}_o^2} = \beta_o - 2\alpha_o\Delta\mathrm{s} + \frac{\Delta\mathrm{s}^2}{\beta_o}\left(1+\alpha_o^2\right)$$

It is clearly a positively defined parabola and we just should find where it has a minimum:

$$\beta'\left(\mathrm{s}^*\right) = 2\left(\mathrm{w}_o + \Delta\mathrm{s}^*\mathrm{w}'_o\right)\mathrm{w}'_o + 2\frac{\Delta\mathrm{s}^*}{\mathrm{w}_o^2} = 0 \to \Delta\mathrm{s}^* = -\mathrm{w}_o^2\frac{\mathrm{w}_o\mathrm{w}'_o}{1+\left(\mathrm{w}_o\mathrm{w}'_o\right)^2} = \frac{\alpha_o\beta_o}{1+\alpha_o^2}$$

$$\beta^* = \beta\left(\mathrm{s}^*\right) = \frac{\beta_o}{1+\alpha_o^2}; \mathrm{w}^* = \sqrt{\frac{\beta_o}{1+\alpha_o^2}}; \mathrm{w}'^* = 0;$$

Now we need just to apply (2) again with $s_o = s^*$:

$$\beta(\mathrm{s}) = \mathrm{w}^2(\mathrm{s}) = \left|\mathrm{w}^* + \frac{i\left(\mathrm{s}-\mathrm{s}^*\right)}{\mathrm{w}^*}\right|^2 = \beta^* + \frac{\left(\mathrm{s}-\mathrm{s}^*\right)^2}{\beta^*}\#.$$

(b) Using (2) again we have:

$$\mathrm{w}(\mathrm{s})e^{i\psi(s)} = \mathrm{w}_o + (s-s_o)\mathrm{s}\left(\mathrm{w}'_o + \frac{i}{\mathrm{w}_o}\right) = \mathrm{w}^* + i\frac{s-s^*}{\mathrm{w}^*} = \mathrm{w}^*\left(1+i\frac{s-s^*}{\beta^*}\right);$$

$$\psi(s) = \tan^{-1}\left(\frac{s-s^*}{\beta^*}\right) \to \psi(s_2)-\psi(s_1) = \tan^{-1}\left(\frac{s_2-s^*}{\beta^*}\right) - \tan^{-1}\left(\frac{s_1-s^*}{\beta^*}\right).$$

Trajectory:

$$x(z) = a\sqrt{\beta(z)}\cos(\psi(z)+\varphi); \beta(z) = \beta^* + \frac{z^2}{\beta^*}; \tan\psi(z) = \frac{z}{\beta^*};$$

$$x'(z) = a\left(\frac{\beta'(z)}{2\sqrt{\beta(z)}}\cos(\psi(z)+\varphi) - \frac{1}{\sqrt{\beta(z)}}\sin(\psi(z)+\varphi)\right)$$

We should note that:

$$\frac{\beta(z)}{\beta^*} = 1 + \frac{z^2}{\beta^{*2}} = 1+\tan^2\psi = \frac{1}{\cos^2\psi}; \tan\psi(s) = \frac{z}{\beta^*};$$

$$x(z) = a\sqrt{\beta(z)}(\cos\psi\cos\varphi - \sin\psi\sin\varphi) = \frac{a\sqrt{\beta^*}}{\cos\psi}(\cos\psi\cos\varphi - \sin\psi\sin\varphi)$$

$$x(z) = a\sqrt{\beta^*}(\cos\varphi - \tan\psi\sin\varphi) = a\sqrt{\beta^*}\left(\cos\varphi - \frac{z}{\beta^*}\sin\varphi\right)$$

e.g. the trajectory is a straight line with constant

$$x'(z) = -\frac{a}{\sqrt{\beta^*}}\sin\varphi$$

(c) Assuming an very long drift

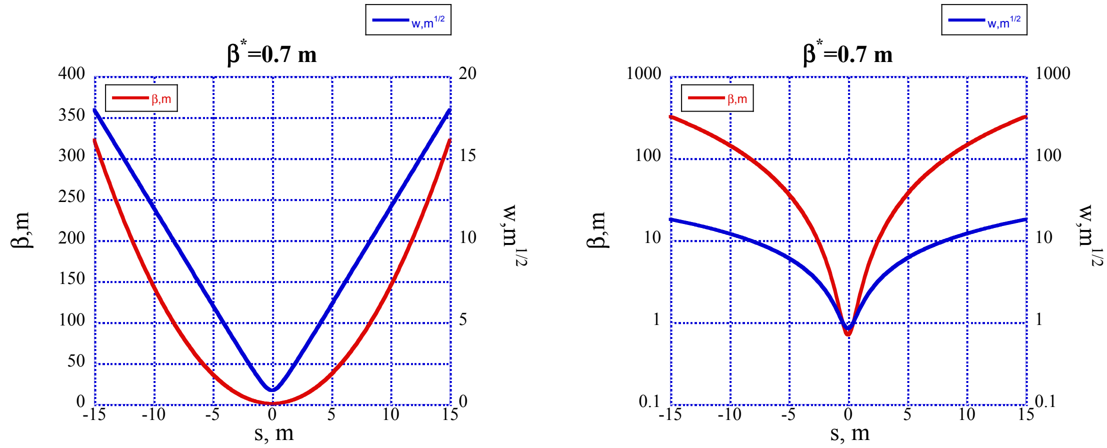
$$s_1 \to -\infty; s_2 \to +\infty$$

$$\psi(s_2)-\psi(s_1) \to \tan^{-1}\left(\frac{\to +\infty}{\beta^*}\right) - \tan^{-1}\left(\frac{\to -\infty}{\beta^*}\right) = \pi$$

Naturally, you can get exactly the same result by integrating the phase advance using

$$\frac{d\psi}{ds} = \frac{1}{\beta(s)} \to \psi(s_2)-\psi(s_1) = \int_{s_1}^{s_2}\frac{ds}{\beta^* + \frac{(\mathrm{s}\text{-}\mathrm{s}^*)^2}{\beta^*}} =$$

$$\tan^{-1}\left(\frac{s_2-s^*}{\beta^*}\right) - \tan^{-1}\left(\frac{s_1-s^*}{\beta^*}\right)$$



Plot of beta-function and beam envelope in 30-m long straight section with β*=0.7 m – typical for RHIC interection region.