

Transverse (Betatron) Motion

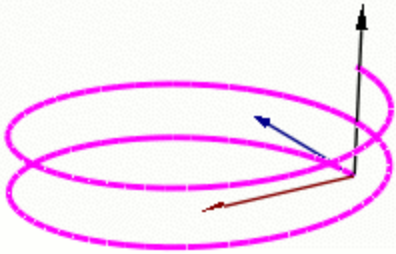
Linear betatron motion

Dispersion function of off momentum particle

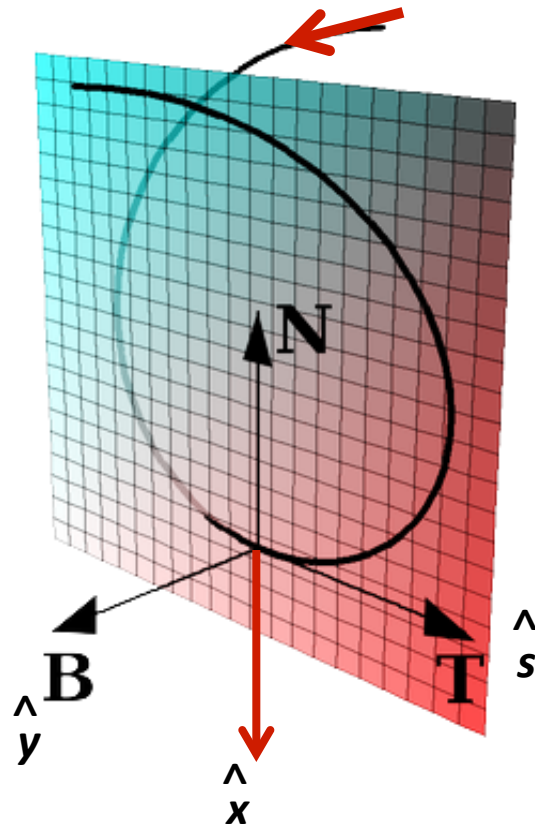
Simple Lattice design considerations

Nonlinearities

Frenet-Serret coordinate system:



1. The tangential vector points toward the direction of beam motion
2. We define the normal vector points outward from the curve. Magnetic field direction y . So that $s = -x \times y$ (electron for this case)



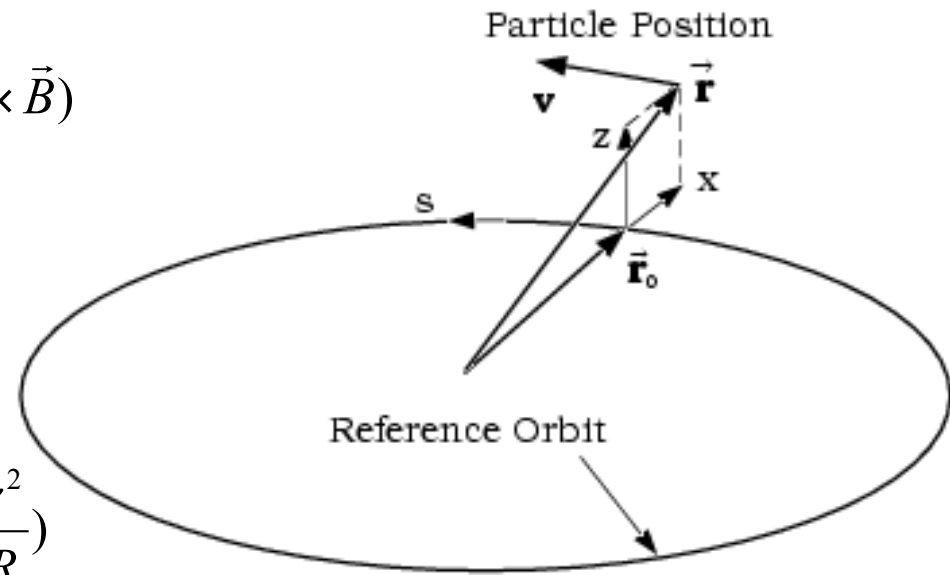
Lorentz force:

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$$

$$F_x = \frac{d}{dt}(\gamma m \dot{x}) = \frac{\gamma m \beta^2 c^2}{R} = -q \beta c B_y$$

$$F_y = \frac{d}{dt}(\gamma m \dot{y}) = q \beta c B_x$$

$$\frac{d}{dt}(\gamma m \dot{x}) \approx \gamma m(\ddot{x} - \frac{v^2}{R}) = \gamma m(\frac{\beta^2 c^2}{R^2} \frac{d^2 x}{d\theta^2} - \frac{v^2}{R})$$



How to transform from the original coordinate system into the Frenet-Serret coordinate system?

$$\theta = \frac{s}{R} = \frac{\beta c t}{R}$$

$$n = -\frac{\rho}{B_0} \left(\frac{\partial B_y}{\partial x} \right)_{x=0} \rightarrow \text{Constant when } B_y \sim x^0 \text{ or } x^1$$

$$\frac{d^2 x}{d\theta^2} + \left(1 + \frac{\rho}{B_0} \frac{\partial B_y}{\partial x}\right) x = 0$$

$$\frac{d^2 y}{d\theta^2} + \left(-\frac{\rho}{B_0} \frac{\partial B_y}{\partial x}\right) y \left(1 + 2 \frac{x}{\rho}\right) = 0$$

$$\frac{d^2 x}{d\theta^2} + (1 - n)x = 0$$

$$\frac{d^2 y}{d\theta^2} + ny = 0$$

$$\frac{dx}{d\theta} = \sqrt{1-n}(-A \sin \sqrt{1-n}\theta + B \cos \sqrt{1-n}\theta) \quad \frac{dy}{d\theta} = \sqrt{n}(-C \sin \sqrt{n}\theta + D \cos \sqrt{n}\theta)$$

Number of transverse oscillation in one beam revolution $\sim \sqrt{1-n} / \sqrt{n}$, weak focusing!

For weak focusing, stable solution requires $0 < n < 1$, which makes transverse (betatron) oscillation having a tune (number of circles per revolution)

$Q_x, Q_y = \sqrt{1-n}, \sqrt{n}$ less than 1. The beam sizes in such machines scales with $C^{1/2} / (1-n)^{1/4}, C^{1/2} / n^{1/4}$

For a modern machine, especially light sources where smaller beam sizes are required to achieve higher beam brightness, strong focusing is required => external focusing magnets (quadrupoles) are often used.

$$x'' + K_x(s)x = \pm \frac{\Delta B_z}{B\rho}, \quad y'' + K_y(s)y = \mp \frac{\Delta B_x}{B\rho}$$

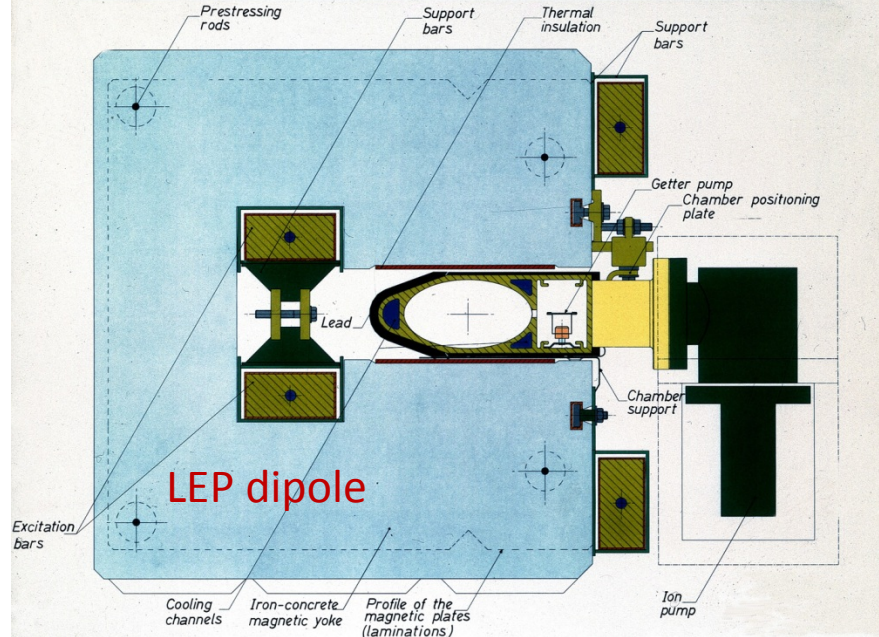
$$K_x(s) = \frac{1}{\rho^2} \mp \frac{B_1}{B\rho}, \quad K_y(s) = \pm \frac{B_1}{B\rho}$$

Natural focusing
from dipoles

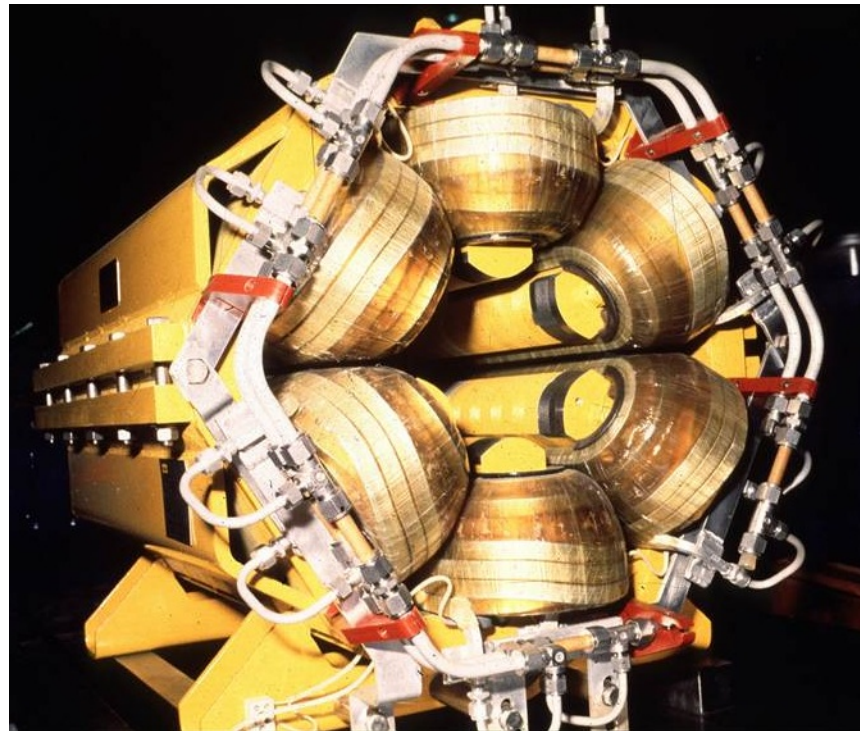
Focusing from
quadrupoles

Higher order
magnet, usually field
errors

CROSS SECTION OF THE DIPOLE MAGNET WITH THE VACUUM CHAMBER

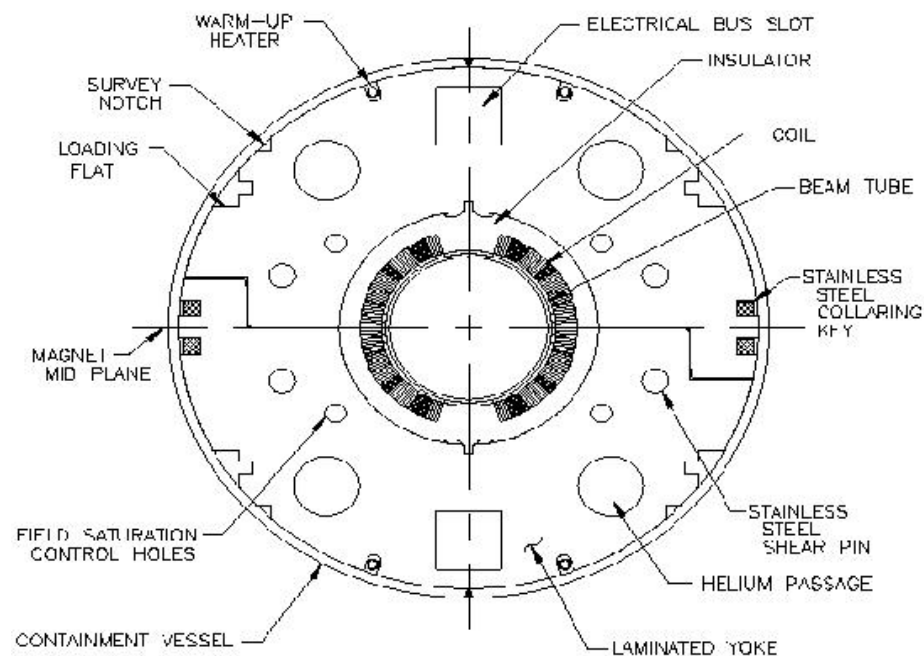


ISR
quadrupole

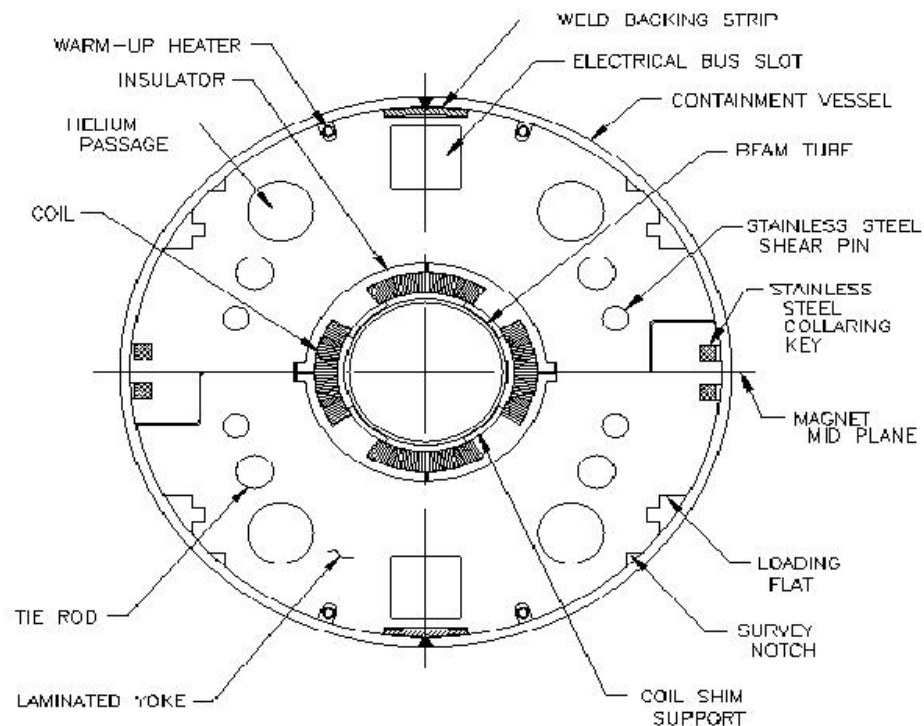


LEP Sextupole

RHIC ARC DIPOLE



RHIC ARC QUADRUPOLE



For two dimensional magnetic field, one can expand the magnetic field using **Beth representation**.

$$\vec{B} = B_x(x, y)\hat{x} + B_y(x, y)\hat{y}$$

$$B_x = -\frac{1}{h_s} \frac{\partial(h_s A_2)}{\partial y} = -\frac{1}{h_s} \frac{\partial A_s}{\partial y}, B_y = \frac{1}{h_s} \frac{\partial(h_s A_2)}{\partial x} = \frac{1}{h_s} \frac{\partial A_s}{\partial x}$$

For $h_s=1$ or $\rho=\infty$, one obtains the multipole expansion:

$$B_y + jB_x = B_0 \sum_n (b_n + ja_n)(x + jy)^n, \quad A_s = \text{Re} \left\{ B_0 \sum_n \frac{1}{n+1} (b_n + ja_n)(x + jy)^{n+1} \right\}$$

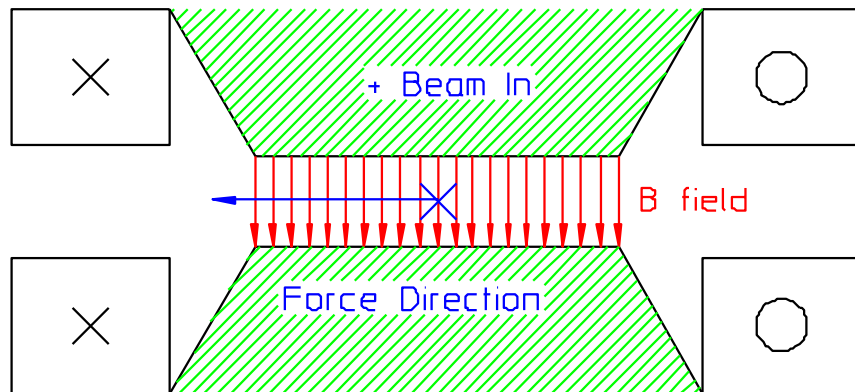
b_0 : dipole, a_0 : skew (vertical) dipole; $B_y = B_0 b_0$, $B_x = B_0 a_0$,

b_1 : quad, a_1 : skew quad; $B_y = B_0 b_1 x$, $B_x = B_0 b_1 y$, $B_y = -B_0 a_1 y$, $B_x = B_0 a_1 x$,

b_2 : sextupole, a_2 : skew sextupole;

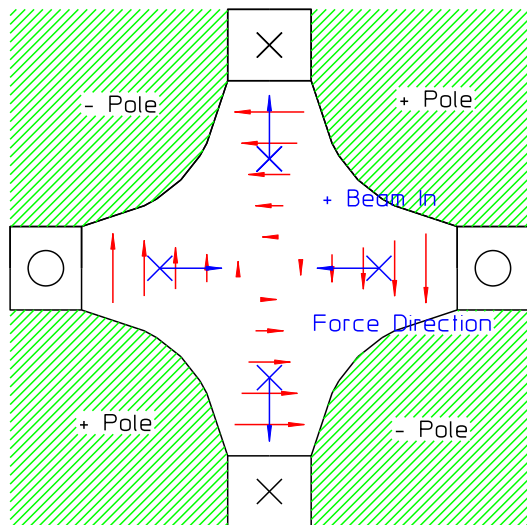
$$\frac{1}{B\rho} (B_y + jB_x) = \mp \frac{1}{\rho} \sum_n (b_n + ja_n)(x + jy)^n$$

+ Current In Positive Pole + Current Out



Negative Pole

+ Current In

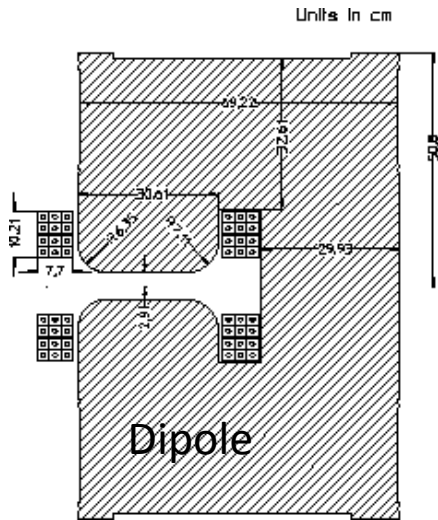


+ Current Out

$$B_0 = \frac{\mu_0 NI}{g}, \quad A_s = B_0 x$$

$$B_1 = \frac{2\mu_0 NI}{a^2}, \quad A_s = \frac{B_1}{2} (x^2 - z^2)$$

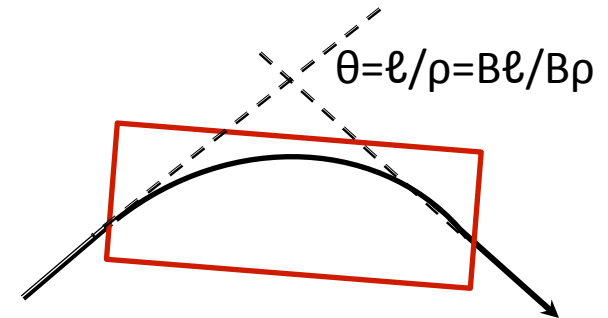
We will learn how to optimize the arrangement of beam control elements in achieving the wanted beam properties.



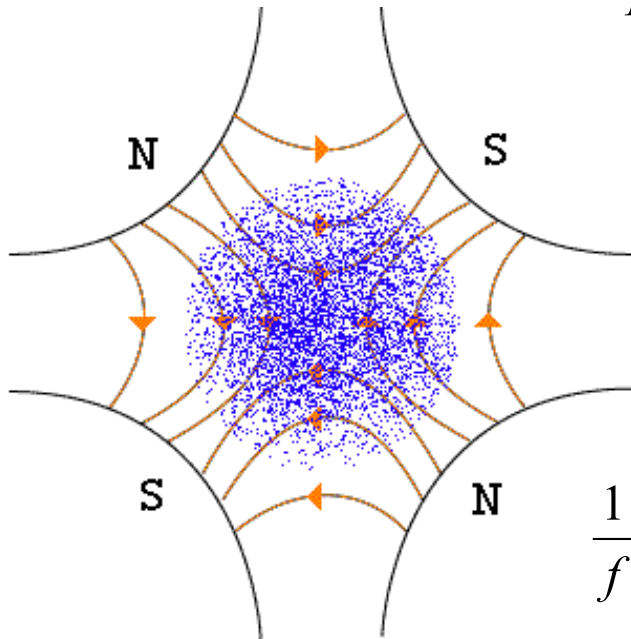
$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$$

$$\gamma m \frac{v^2}{\rho} = q v B$$

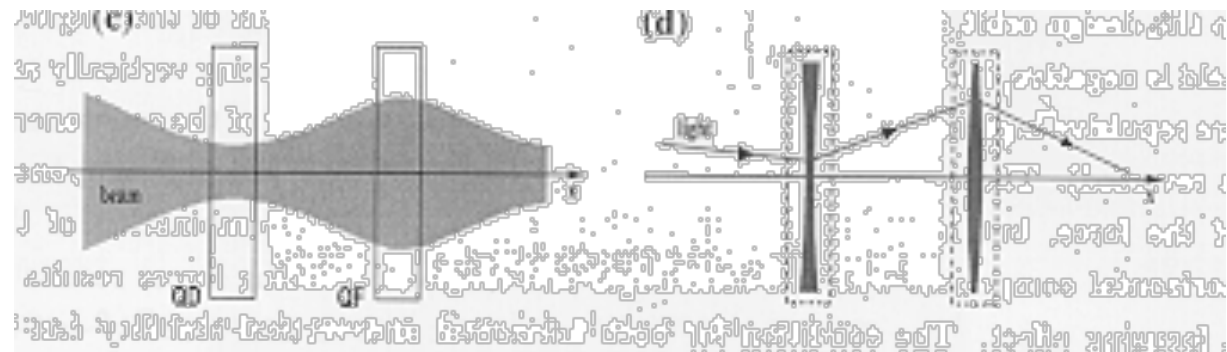
$$\rho = \frac{\gamma m v}{q B} = \frac{p}{q B}$$



quadrupole



$$B \rho [T \cdot m] = \frac{p}{q} = \frac{A}{Z} \times 3.33564 \times p [GeV / c / u]$$



$$\frac{1}{f} = \mp \frac{B_1 l}{B \rho}$$

$f > 0$, if focusing, $f < 0$ if defocusing

How to solve the Hill's equation?

$$x'' + K_x(s)x = \frac{\Delta B_y}{B\rho}, \quad y'' + K_y(s)y = -\frac{\Delta B_x}{B\rho}$$

Ideal Linear accelerator:

$$x'' + K_x(s)x = 0, \quad y'' + K_y(s)y = 0$$

Let X represent x or y :

$$X'' + K(s)X = 0,$$

$$X(s) = \begin{pmatrix} X(s) \\ X'(s) \end{pmatrix} = M(s, s_0) \begin{pmatrix} X(s_0) \\ X'(s_0) \end{pmatrix}$$

$M(s, s_0)$ is the betatron transfer matrix. For any two linearly independent solutions y_1, y_2 of Hill's equation, the Wronskian is independent of time

$$W(y_1, y_2, s) \equiv y_1 y_2' - y_1' y_2, \quad \frac{dW}{ds} = 0.$$

$$W(s) = [\det M] W(s_0)$$

$\text{Det}(M(s_2, s_1)) = 1$

